

# Time - Frequency Analysis

## Linear Analysis

Decompose the signal  $x(t)$  into elementary "time-frequency atoms" (e.g., Gabor functions) :

$$x(t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S_x(u, f) g(t-u) e^{j2\pi f t} df du$$

"weights"      Gabor functions

The weights in this superposition reflect the distribution of signal energy in time and freq.

## Quadratic Analysis

Goal is to find a time-freq energy density  $P_x(t, f)$  satisfying

$$\int_{-\infty}^{\infty} P_x(u, f) du = |X(f)|^2$$

$$\int_{-\infty}^{\infty} P_x(t, v) dv = |X(t)|^2$$

Clearly, this requires that  $P_x$  is a quadratic function of  $X(t)$ .

## Linear Time-frequency Analysis

The basic idea is to represent a signal  $x$  in terms of elementary "building-block" signals that are well concentrated in time and frequency. These elementary signals are called "time-frequency atoms".

Let  $\{\phi_\gamma\}_{\gamma \in \Gamma}$  be a family of TF atoms. Here, the index  $\gamma$  may indicate both time- and frequency localization; e.g.,

$$\gamma = (t_0, f_0)$$

Let us assume that all atoms are normalized

$$\|\phi_\gamma\|_2 = \left( \int |\phi_\gamma(t)|^2 dt \right)^{1/2} = 1.$$

Then we can define a linear time-frequency operator  $T$  as

$$Tx(\gamma) = \int_{-\infty}^{\infty} x(t) \phi_{\gamma}^*(t) dt \\ = \langle x, \phi_{\gamma} \rangle$$

Also, by Parseval's Theorem

$$Tx(\gamma) = \int_{-\infty}^{\infty} X(f) \Phi_{\gamma}^*(f) df$$

where  $X, \Phi$  are Fourier transforms of  $x, \phi$ .

So, if  $\phi_{\gamma}$  is concentrated in a neighborhood about time  $t_0$ , then

$Tx(\gamma)$  depends only on values of  $x$  in this neighborhood. Similarly, if

$\Phi_{\gamma}$  is concentrated in a neighborhood about frequency  $f_0$ , then  $Tx(\gamma)$  only depends of  $X(f)$  in this neighborhood.

Ex.

► Windowed Fourier Transform

$$\phi_r(t) = e^{j2\pi f_0 t} w(t - t_0)$$

$w(t - t_0)$   
window function

Wavelet Transform

$$\phi_r(t) = \frac{1}{\sqrt{s_0}} \psi\left(\frac{t - t_0}{s_0}\right)$$

►  $\gamma = (t_0, s_0)$  ,  $\psi(t)$  = "mother" wavelet  
     ↗ scale instead  
        of frequency

$$\text{scale } s_0 \propto \frac{1}{f_0}$$

smaller scale  
 $\iff$  higher frequency

## Time-Frequency Concentration

The energy of  $\Phi_r$  is centered at

$$t_r = \int_{-\infty}^{\infty} t |\Phi_r(t)|^2 dt$$

in time, and in frequency at

$$f_r = \int_{-\infty}^{\infty} f |\Phi_r(f)|^2 df$$

The energy "spread" about  $(t_r, f_r)$  is measured by

$$\sigma_t^2(r) = \int_{-\infty}^{\infty} (t - t_r)^2 |\Phi_r(t)|^2 dt$$

$$\sigma_f^2(r) = \int_{-\infty}^{\infty} (f - f_r)^2 |\Phi_r(f)|^2 df$$

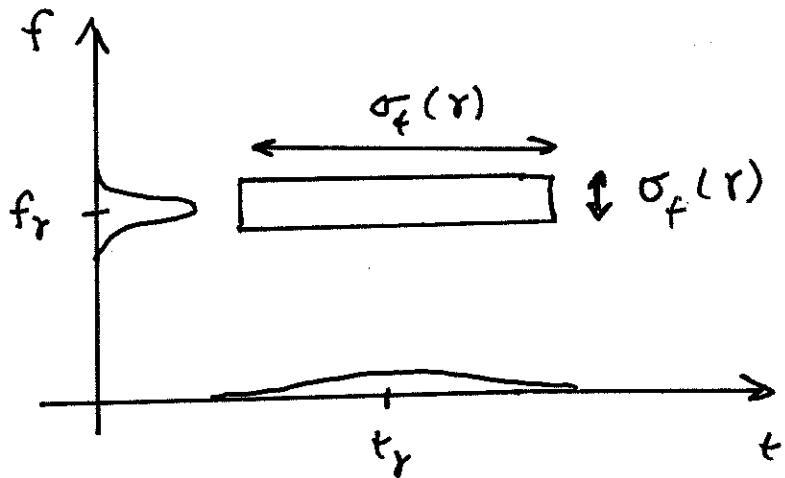
Heisenberg Uncertainty

$$\Rightarrow \sigma_t(r) \sigma_f(r) \geq$$

## Time-Frequency Energy Density

$$|\langle x, \phi_r \rangle|^2 = \left| \int_{-\infty}^{\infty} x(t) \phi_r^*(t) dt \right|^2$$

measures the energy of  $x$   
in the neighborhood of  $(t_r, f_r)$ .



$|Tx(\gamma)|^2$  is a time-frequency  
energy distribution function of  $x$ .

# Short-Time Fourier Transform (STFT)

Replace global Fourier analysis

$$X(f) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi ft} dt$$

with a sequence of local analyses

with respect to a moving observation  
window

$$\begin{aligned} Sx(t, f) &= \int_{-\infty}^{\infty} x(u) e^{-j2\pi fu} w(u-t) du \\ &\quad \underbrace{\qquad\qquad\qquad}_{\text{window about time } t.} \\ &= \int_{-\infty}^{\infty} X(v) W(v-f) e^{-j2\pi(v-f)t} dv \\ &\quad \underbrace{\qquad\qquad\qquad}_{\text{window about frequency } f.} \end{aligned}$$

## Common windows

Gaussian  $\Rightarrow$  Gabor analysis

rectangle, Hamming, Blackman

$w(t)$  assumed to be real, symmetric  
and  $\int |w(t)|^2 dt = 1$ .

$$\underline{\text{Ex.}} \quad x(t) = e^{j2\pi f_0 t}$$

$$\begin{aligned} S_x(t, f) &= \int_{-\infty}^{\infty} x(u) e^{-j2\pi f u} w(u-t) du \\ &= \int_{-\infty}^{\infty} X(v) w(v-f) e^{-j2\pi(v-f)+} dv \end{aligned}$$

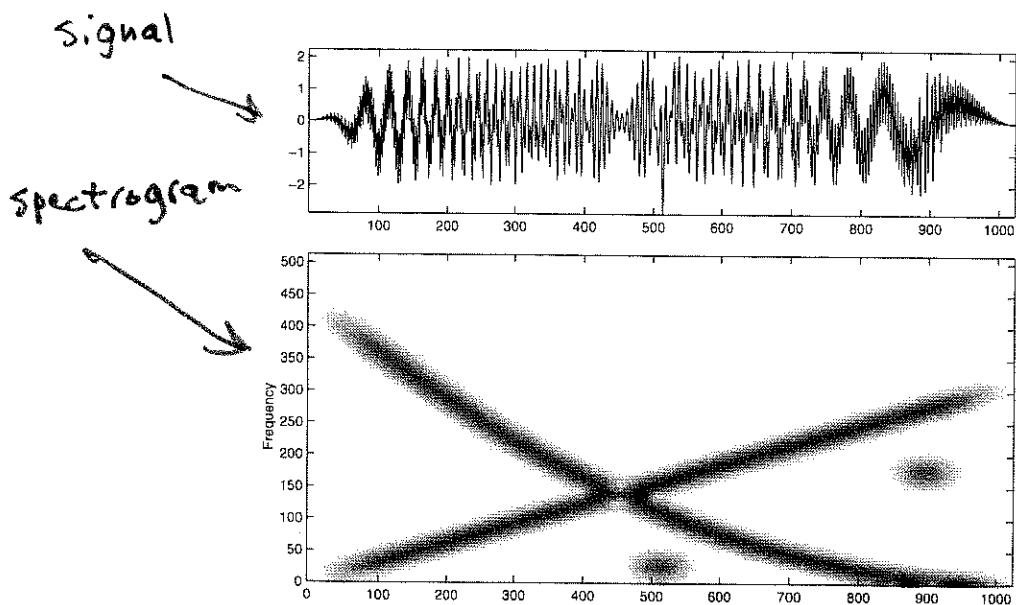
$$\underline{\text{Ex.}} \quad x(t) = \delta(t - t_0)$$

## Spectrogram

The energy distribution resulting from the STFT is called the spectrogram

$$\begin{aligned} P_{S_x}(t, f) &= |S_x(t, f)|^2 \\ &= \left| \int_{-\infty}^{\infty} x(u) e^{-j 2\pi f u} w(u-t) du \right|^2 \end{aligned}$$

$$\text{Ex. } w(u) = \frac{1}{(\pi \sigma^2)^{1/4}} e^{-u^2/2\sigma^2}, \quad \sigma = 50$$



## Time-frequency Resolution of Spectrogram

Because  $w(u)$  is assumed to be real and symmetric, both the time and frequency spread about  $(t, f)$  is independent of  $(t, f)$ :

$$\begin{aligned}\sigma_t^2 &= \int_{-\infty}^{\infty} (t-u)^2 |w(t-u) e^{-j2\pi fu}|^2 du \\ &= \int_{-\infty}^{\infty} u^2 |w(u)|^2 du, \text{ indep. of } t\end{aligned}$$

$$\begin{aligned}\sigma_f^2 &= \int_{-\infty}^{\infty} (f-v)^2 |W(f-v) e^{-j2\pi(v-f)t}|^2 dv \\ &= \int_{-\infty}^{\infty} v^2 |W(v)|^2 dv, \text{ indep. of } f\end{aligned}$$

$\Rightarrow$  time resolution  $\sigma_t$  and frequency resolution  $\sigma_f$  are independent of position in the TF plane.